Monte Carlo Calculation for Production Cross-Sections of Projectile’s Isotopes from Therapeutic Carbon and Helium Ion Beams in Different Materials

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Background: Isotopes of the projectile may be produced along the beam path during the irradiation of a target by a heavy ion due to inelastic interactions with the media. This study analyzed the production cross-section of carbon (C) and Helium (He) projectile’s isotopes resulting from the interactions of these beams with different materials along the beam path.

Materials and Methods: In this study, we transport C and He ion beams through different materials. This transportation was made by the Monte Carlo simulation. Particle and Heavy Ion Transport code System (PHITS) has been used for this calculation.

Results and Discussion: It has been found that $^{10}\text{C}$, $^{11}\text{C}$, and $^{13}\text{C}$ from the $^{12}\text{C}$ ion beam and $^{3}\text{He}$ from the $^{4}\text{He}$ ion beam are significant projectile’s isotopes that have higher flux than other isotopes of these projectiles. The $^{4}\text{He}$ ion beam has a higher projectile’s isotope production cross-section along the beam path, which adds more impurities to the beam than the $^{12}\text{C}$ ion beam. These projectile’s isotopes from both the $^{12}\text{C}$ and $^{4}\text{He}$ beams have higher production cross-sections in hydrogenous materials like water or polyethylene.

Conclusion: It is important to distinguish these projectile’s isotopes from the primary beam particles to obtain a precise and accurate cross-section result by minimizing the error during measurement with a nuclear track detector. This study will show the trend of the production probability of projectile’s isotopes for these ion beams.

Keywords: Projectile’s Isotope, Isotope Production Cross-Section, Monte Carlo, Particle and Heavy Ion Transport Code System

Introduction

Ion beam accelerators play a significant role in medical physics [1], radiation physics, accelerator physics, ion implantation, material science, space radiation [2], and cancer therapy [3]. Therefore, it is of paramount importance to understand the mechanism of the interaction between ions and matter to effectively use the ion beams from an accelerator. The total reaction cross-section is one of the most fundamental quantities for understanding the nucleus–nucleus interaction mechanism. However, the total reaction cross-section is difficult to measure experimentally due to various reaction channels. On the other hand, since most reactions undergo a proton change, the total
charge changing cross-section ($\sigma_{\text{TCC}}$) is an alternative and effective method to understand the reaction cross-section [4]. $\sigma_{\text{TCC}}$ measures the probability of changing the charge (proton) of a projectile. It is also important to understand the structure and stability of the nucleus [5] and particle spectrum in a cosmic ray [6]. Many experiments have measured $\sigma_{\text{TCC}}$ in different interactions since 1970 [7–9]. However, neutron changing cross-sections ($^{12}$C to $^3$C,$^{13}$C isotopes) were excluded from this $\sigma_{\text{TCC}}$ measurement.

To conduct this type of experiment, researchers use both active detectors, such as the Cherenkov detector [10], Si detector, different types of scintillators, position sensitive detectors [7, 8, 11], and time projection chambers [12] and passive detectors, essentially track detectors such as CR-39 [13–18] and nuclear emulsion [19]. Track detectors have been used for ion interaction experiments since 1960s [20–22]. Passive detectors such as CR-39 have certain advantages over active detectors in detecting fragments with a range in ‘μm’ scale, multiple fragments from a single interaction, etc. However, the CR-39 detector has a limitation of low mass resolution [23, 24]. As a result, these detectors cannot differentiate between projectiles and their respective isotopes, which is important for experiments such as $\sigma_{\text{TCC}}$ measurement. The incapability of these detectors to distinguish projectiles and isotopes of the projectile is one of the main sources of error in the experiment. Therefore, it is important to determine the projectile isotopes that are produced before coming to the target because of the interacting materials in the beam line and those inside the target for beam purity.

We analyzed $^{12}$C and $^4$He as the main ion beam for cancer therapy in this study. Ion beam contamination due to projectile’s isotopes is one of the major and longstanding experimental issues that must be considered [25], which hinders the measurement. Kozma et al. [26] and Kaki [27] attempted to measure the production cross-section of these projectile’s isotopes for different interactions. This study estimated the production cross-sections of projectile’s isotopes for $^{12}$C and $^4$He ion beams in different media with the Monte Carlo calculation by Particle and Heavy Ion Transport code System (PHITS) [28]. Recently, the use of $^4$He ion has renewed interest for certain therapeutic conditions due to less lateral scattering than the proton and lower impact of the fragmentation tail than in $^{12}$C ion irradiation. A clinical trial of $^4$He ion at the Heidelberg Ion Beam Therapy Center has just started [29]. So, the production cross-section of isotopes from the $^4$He ion beam has also been analyzed in this study.

There are very few experimental data on projectile’s isotope production cross-sections ($\sigma_{\text{ipc}}$) measured for different interactions. Therefore, our calculated data and the analysis of the trend of projectile’s isotope production in different target materials in this study could be useful for researchers in this field.

Materials and Methods

In this study, JAERI Quantum Molecular Dynamics 2.0 (IQMD 2) [30] incorporated in PHITS version 3.22 was used to simulate the transport of the $^{12}$C ion beam, and Liège Intranuclear Cascade model (INCL4.6) [31] has been used for the transportation of $^4$He ion beams inside materials such as scatterer (Ta) with density of 16.6 g/cm$^3$, beam monitors (Al) with density of 2.7 g/cm$^3$, and range shifter (polymethyl methacrylate [PMMA] (C$_3$O$_2$H$_5$)$_n$) with density of 1.91 g/cm$^3$ [32] along the beam path in accelerator facilities such as Heavy Ion Medical Accelerator in Chiba (HIMAC), Japan [33]. Other targets, Al, C with density of 2.26 g/cm$^3$, CR-39 with density of 1.31 g/cm$^3$, water (H$_2$O) with density of 1.0 g/cm$^3$, and polyethylene (CH$_2$) with density of 0.89 g/cm$^3$ were also investigated to understand the tendency of projectile’s $\sigma_{\text{ipc}}$ for both the $^{12}$C and $^4$He ion beams. The target thickness was considered 0.5 mm. The simulation was performed for 1 million particles in each energy for each interaction so that the statistical error of the calculated quantity remains less than 2%. The following Fig. 1A shows the simulation setup for $^{12}$C and PMMA interaction at 135 MeV/n with a target thickness of 0.5 mm, and the beam radius was 5 cm during the simulation. Due to the better visualization, the beam of a 2 cm radius has been visualized in this figure. Fig. 1B shows the travelling of the $^{12}$C ion beam through the PMMA.

1. Calculation of Projectile’s Isotope Production Cross-Section

The track detection mechanism of track detectors, such as CR-39 is based on the incident ion’s linear energy transfer (LET) inside this detector. The LET threshold and LET resolution of a CR-39 (TD-1 type) detector are approximately 5 keV/μm and 1 keV/μm, respectively. Fig. 2A shows the LET as a function of energy for $^{11}$C, $^{12}$C, and $^{13}$C in CR-39. It can be observed in the Fig. 2A that the difference in LET among these isotopes is insignificant. As shown in Fig. 2B, the average difference in LET between $^{12}$C and $^{13}$C ($\text{LET}_{\text{c-12}} - \text{LET}_{\text{c-13}}$) and that between $^{12}$C and $^{14}$C ($\text{LET}_{\text{c-12}} - \text{LET}_{\text{c-14}}$) is approxi-
The production of C and He isotopes from their projectiles is as follows [34]:

\[
\sigma_{\text{ipc}} = \frac{M_p}{\rho N_A x} \ln \left( \frac{N_{\text{in}} - N_{\text{ne}}}{N_{\text{in}}} \right)
\]

where, \( N_i \) is the Avogadro number, \( N_{\text{in}} \) is the number of incoming particles, and \( N_{\text{ne}} \) is the number of interactions in which an isotope has been produced by the change of neutron(s). \( M_p, \rho, \) and \( x \) are the molar mass, density, and thickness of the target, respectively. Thus, the production cross-section of a specific isotope (e.g., \(^{11}\)C) is

\[
\sigma_{C-11} = \frac{M_p}{\rho N_A x} \ln \left( \frac{N_{\text{in}} - N_{C-11}}{N_{\text{in}}} \right)
\]

where, \( N_{C-11} \) is the total number of interactions in which the mass of the projectile changed to that of \(^{11}\)C by the removal of neutron. Thus, the total isotope production cross-section (\( \sigma_{\text{tot}} \)) is the sum of the production cross-sections of all types.
of isotopes. That is,

$$\sigma_{\text{TPC}} = \sum_{n=0}^{15} \sigma_{C-n}$$  \hspace{1cm} (3)$$

where, \(n\) is the mass number, which defines the specific isotopes. We cannot distinguish the isotopes of a projectile by experimental measurement using a track detector because of the low mass resolution. Consequently, the total production cross-section of projectile isotopes is likely to be measured simultaneously during the measurement of the charge changing cross-section, and it is equivalent to the mass changing cross-section.

Thus, we calculated the production cross-sections of isotopes of projectiles using PHITS by identifying the event-by-event interactions of projectiles with the target while changing their neutron \(N_n\).

Results and Discussion

1. Flux of Projectile’s Isotopes along Beam Path

The transport of \(^{12}\)C and \(^{4}\)He ion beams of 135, 290, and 400 MeV/n having a diameter of 5 cm through the materials (Ta, Al, and PMMA) of 0.5 mm thickness along the beam path has been simulated. We applied PHITS for this simulation to assess the beam contamination by the projectile’s isotopes produced in the beam path. These energies are in the therapeutic region provided by HIMAC. Fig. 3A and 3B show the flux of isotopes as a function of energy in the PMMA for \(^{12}\)C ion beams of 135 and 400 MeV/n, respectively. The figures reveal that \(^{4}\)He, \(^{12}\)C, \(^{13}\)C, and \(^{15}\)N have significantly low energy with the least flux. Hence, these are too insignificant to be considered because it is highly challenging to reach the

Fig. 3. (A) The flux of isotopes as a function of energy in the PMMA along the beam path for the \(^{12}\)C ion beam of 135 MeV/n. (B) The same for the \(^{12}\)C ion beam of 400 MeV/n. (C, D) The same for \(^{4}\)He and PMMA interactions, from where it is seen that \(^{4}\)He and \(^{14}\)He are the significant isotopes for the \(^{4}\)He ion beam. PMMA, polymethyl methacrylate.
target. Only the $^{10}$C, $^{11}$C, and $^{13}$C isotopes contribute significantly through their higher flux and energy for the contamination of the beam and results of cross-section measurements. It can also be observed that the production of these C isotopes is higher at lower energy. In Fig. 3, we can see that the production of isotopes is higher at 135 MeV/n than 400 MeV/n in the same target, PMMA. This may substantially influence the beam purity before the target is reached. It is also found that $^4$He and $^6$He are the significant isotopes from the $^4$He ion beam that are produced along the beam path for the same interacting materials at the same energies, as shown in Fig. 3C and 3D.

2. Projectile’s Isotope Production Cross-Sections

1) Carbon projectile isotopes

Each interaction of the projectile from which the projectile’s isotopes ($^{12}$C–$^{15}$C) can be produced was identified separately using the [T-userdefined] tally of PHITS. Then, the production cross-section of specific projectile’s isotopes was calculated from the number of specific interactions using Equation (2). Table 1 shows the $\sigma_{ipc}$ for specific projectile’s isotopes ($^{12}$C–$^{15}$C) for different energies in the materials along the beam path of the accelerator system. It can be observed that $^{13}$C and $^{14}$C have the maximum production probabilities in each material. The value enclosed in parentheses shows the uncertainty calculated according to the error propagation.

The production probabilities of different isotopes in certain target materials, such as Al, C, CR-39, H$_2$O, and CH$_3$, were calculated in this study to understand the trend, and these are shown in Fig. 4. $^{13}$C, $^{14}$C, and $^{15}$C are the most producible isotopes in all the target materials, as observed from these figures and Table 1, which represent the isotopes of the projectile in the materials along the beam path. It is observed that $^{12}$C, $^{13}$C, and $^{14}$C have a significantly higher cross-section in PMMA, H$_2$O, and CH$_3$ than other targets. The production cross-section of $^{14}$C was significantly higher than that of $^{13}$C and $^{15}$C in all the targets.

The calculated production cross-sections of $^{13}$C and $^{14}$C

Table 1. Production Cross-Section (mb) of C Isotopes ($^{12}$C–$^{15}$C) for Different Energies of $^{12}$C Ion Beam in Different Materials along the Beam Path

<table>
<thead>
<tr>
<th>Materials</th>
<th>Ta</th>
<th>Energy (MeV/n)</th>
<th></th>
<th></th>
<th></th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td>135</td>
<td>290</td>
<td>400</td>
<td>135</td>
</tr>
<tr>
<td>Production cross-section (mb)</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$\sigma_{12C}$</td>
<td>2.2</td>
<td>0.4</td>
<td>0.4</td>
<td>0.3</td>
<td>1.3</td>
</tr>
<tr>
<td>$\sigma_{13C}$</td>
<td>2.0</td>
<td>2.0</td>
<td>3.0</td>
<td>8.0</td>
<td>5.0</td>
</tr>
<tr>
<td>$\sigma_{14C}$</td>
<td>45.0</td>
<td>39.0</td>
<td>35.0</td>
<td>70.0</td>
<td>59.0</td>
</tr>
<tr>
<td>$\sigma_{15C}$</td>
<td>-</td>
<td>0.4</td>
<td>-</td>
<td>31.2</td>
<td>16.9</td>
</tr>
<tr>
<td>$\sigma_{ipc}$</td>
<td>-</td>
<td>-</td>
<td>-</td>
<td>3.0</td>
<td>4.3</td>
</tr>
<tr>
<td>We can observe that the maximum production probabilities are for $^{13}$C and $^{15}$C in each material. The value in parentheses indicates the uncertainty. PMMA, polymethyl methacrylate; $\sigma_{ipc}$, total isotope production cross-section.</td>
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</tbody>
</table>

Fig. 4. The production cross-sections of $^{13}$C, $^{14}$C, and $^{15}$C in different target materials are shown in (A, B, C), respectively. The descending order of production cross-section of $^{13}$C in different targets is H$_2$O > CH$_3$ > C > Al > CR-39. The same for $^{14}$C is H$_2$O > CH$_3$ > C > CR-39 > Al. The same for $^{15}$C is H$_2$O > Al > CR-39 > C > CH$_3$. Carbon-11 ($^{11}$C) has a higher production cross-section than $^{10}$C in all materials.
were compared with the experimental results, as shown in Fig. 5. The experimental results were collected from the references [35–39]. In Fig. 5, the circles and squares indicate the experimental production cross-sections of $^{12}$C and $^{11}$C, respectively, for $^{12}$C and C interaction. The blue and orange colored dashed lines indicate the production cross-sections calculated by PHITS for $^{12}$C and $^{11}$C, respectively, for the same interaction. The calculation process and equation can reproduce the trend of the experimental results within the acceptable range, as observed in Fig. 5.

2) Helium projectile isotopes

In this study, the production cross-section of the projectile’s isotopes for a $^{4}$He ion beam has also been analyzed for the aforementioned target materials as well as for the material along the beam path. During the interaction with the target materials, the produced projectile isotopes from the $^{4}$He ion beam are $^{3}$He and $^{4}$He. The following Table 2 shows the production cross-sections of these He isotopes during the interaction with material along the beam path at the nominal energies 135, 290, and 400 MeV/n through the materials Ta, Al, and PMMA.

Fig. 6A shows the production cross-sections of He isotope in other target materials. This figure shows the produced He isotopes during the interaction of $^{4}$He and C at 80 MeV/n. $^{3}$He is the most significant projectile isotope that is produced and can play an important role during the measurement of charge changing cross-section with track detectors. The production cross-section of $^{3}$He isotopes has been calculated and compared for different targets, as shown in Fig. 6B. It is found that for both H$_2$O and CH$_3$, $^{3}$He has the highest production cross-section.

**Conclusion**

Accelerator-based ion beam has an important role in many branches of physics, including heavy ion cancer therapy, the study of heavy ion interaction, etc. When the beam of a projectile is going through the material before reaching the target, the projectile’s isotopes could be produced, which may contaminate the beam due to the material along the beam path like Ta, Al, range shifter (PMMA), etc. These projectile’s isotopes are one of the main sources of error during the cross-section measurement by the nuclear track detector. There are very few experimental data for the production cross-section of the projectile’s isotope yet. Therefore, we tried to observe the trend of the production cross-section of projectile isotopes for the $^{12}$C and $^{4}$He ion beams in the materials

![Fig. 5. A comparison between the calculated and experimental production cross-sections of $^{12}$C and $^{11}$C for the $^{12}$C and C interactions. The experimental results have been collected [35–39]. The circles and squares indicate the production cross-sections of $^{12}$C and $^{11}$C, respectively. The blue and orange colored dashed lines denote the calculated production cross-sections for $^{12}$C and $^{11}$C, respectively. It is observed that the experimental results are in agreement with the calculated results. PHITS, Particle and Heavy Ion Transport code System.](image-url)

**Table 2. Production Cross-Section (mb) of He Projectile’s Isotopes from $^{4}$He Ion Beam in Different Materials along the Beam Path for Different Energies**

<table>
<thead>
<tr>
<th>Materials</th>
<th>Energy (MeV/n)</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>Ta</td>
</tr>
<tr>
<td>Production cross-section (mb)</td>
<td></td>
</tr>
<tr>
<td>$\sigma_{12c-2}$</td>
<td>71.3</td>
</tr>
<tr>
<td>$\sigma_{12c-6}$</td>
<td>7.0</td>
</tr>
<tr>
<td>$\sigma_{12c}$</td>
<td>78.3 (1.1)</td>
</tr>
</tbody>
</table>

The maximum production probability is for $^{3}$He in each material. The value in parentheses indicates the uncertainty. PMMA, polymethyl methacrylate; $\sigma_{12c}$, total projectile isotope production cross-section.
(Ta, Al, and PMMA) along the beam path for three nominal beam energies, which are 135, 290, and 400 MeV/n, since these energies are in the therapeutic range provided by the HIMAC. The calculations were performed by the Monte Carlo calculation code, PHITS. It is found that $^3$He is the most producible projectile’s isotope from the $^4$He ion beam, whereas $^{10}$C, $^{11}$C, and $^{13}$C are the most producible projectile’s isotope from the $^{12}$C ion beam in these energies. Among the beam materials, Al, Ta, and PMMA, along the beam path, $^3$He have higher production cross-sections in Al and Ta, whereas $^{10}$C, $^{11}$C, and $^{13}$C have higher production cross-sections in PMMA.

Except for the materials along the beam path, other materials, C, H$_2$O, CH$_2$, and CR-39, were considered because of their tissue-like properties to observe the tendency of the projectile’s $\sigma_{ipc}$. It is found that the projectile’s $\sigma_{ipc}$ is higher in hydrogenous materials like H$_2$O and CH$_2$. It can be concluded that the isotopes $^{10}$C, $^{11}$C, and $^{13}$C from the $^{12}$C ion beam and $^3$He from the $^4$He ion beam play an important role in hindering the experimental measurement of the total charge changing cross-section in the Al, C, CR-39, H$_2$O, and CH$_2$ targets. These isotopes were higher in flux at lower energies. Carbon-11 ($^{11}$C) has a much higher production cross-section among the three C projectile isotopes in all the target materials mentioned above.

As determined in this study, the role of these projectile’s isotopes may explain why the experimental result is higher than the calculated result for the $\sigma_{TCC}$ particularly in the lower energy region. Thus, this issue should be addressed effectively to obtain a precise result with minimum uncertainties in an experiment involving an ion beam from an accelerator. The observations of this analysis can be used to minimize the error generated during experiments on heavy ion interactions with track detectors.

**Conflict of Interest**

No potential conflict of interest relevant to this article was reported.

**Ethical Statement**

This article does not contain any studies with human participants or animals performed by any of the authors.

**Data Availability**

The data that support the findings of this study are available on request from the corresponding author (Dr. Quazi Muhammad Rashed Nizam).

**Author Contribution**

Conceptualization: Rashed Nizam QM. Formal analysis: all authors. Writing - original draft: Ahmed A, Ahmed I. Writing - review and editing: Rashed Nizam QM. Approval of final manuscript: all authors.
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